

Entanglement Witnesses



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Separable States

Entangled States



SEP

$$\rho^{sep} = \sum_i P_i \rho_i \otimes \sigma_i$$

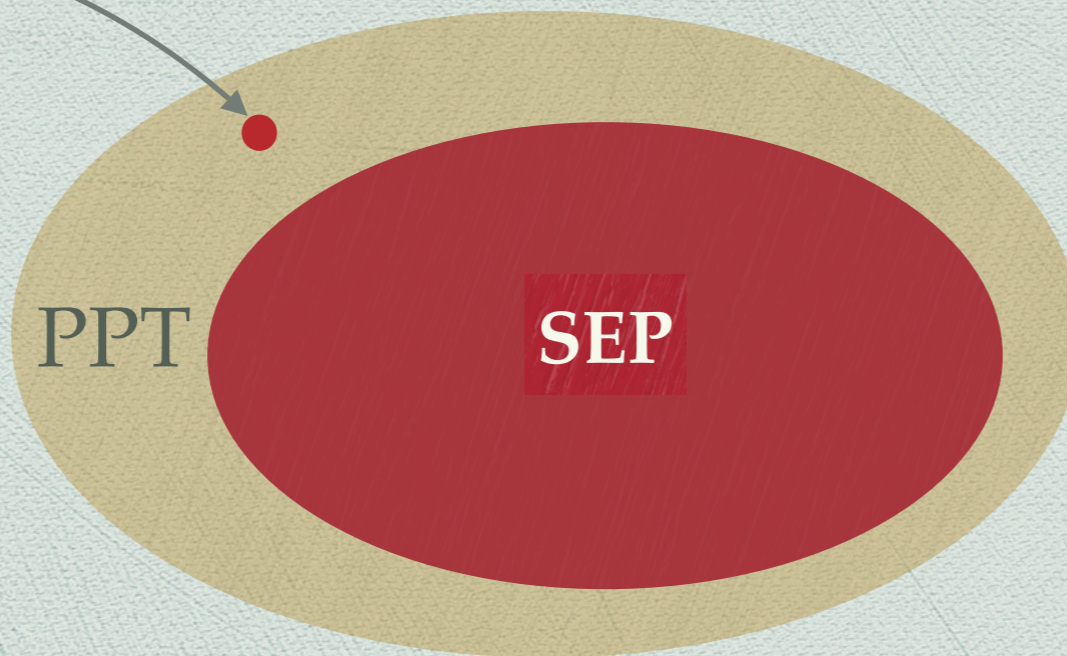
Is this state separable?

$$\rho = \frac{1}{4} \begin{pmatrix} 1-p & 0 & 0 & 0 \\ 0 & p+1 & -2p & 0 \\ 0 & -2p & p+1 & 0 \\ 0 & 0 & 0 & 1-p \end{pmatrix}$$

Positive Partial Transpose (PPT)

$$\rho = \sum_i P_i \rho_i \otimes \sigma_i \quad \longrightarrow \quad \rho^\Gamma \equiv (I \otimes T)\rho = \sum_i P_i \rho_i \otimes \sigma_i^T$$

Bound
entangled state

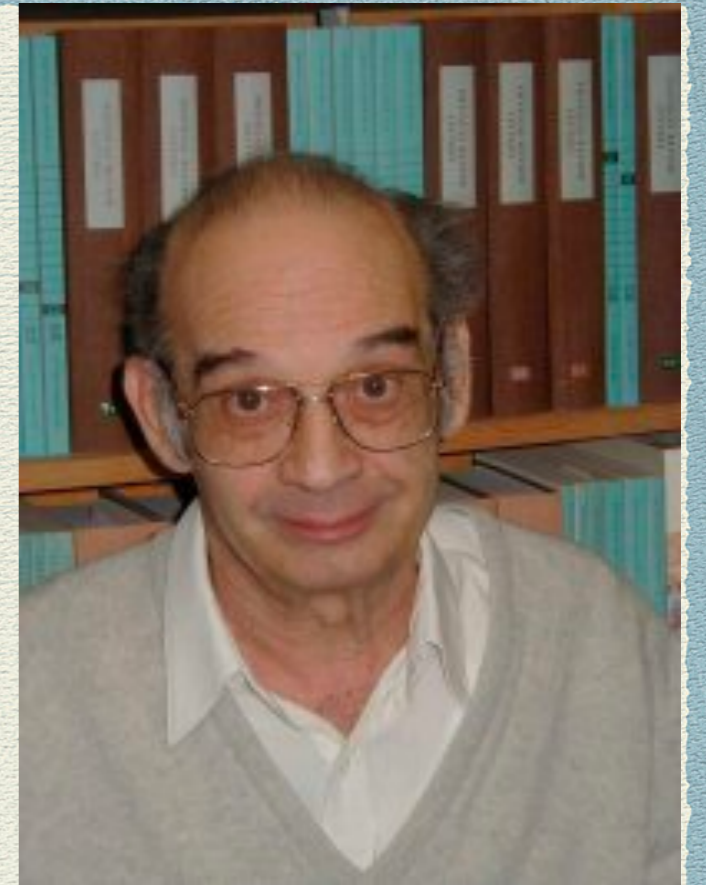


Peres Criteria

$$\rho = \frac{1}{4} \begin{pmatrix} 1-p & 0 & 0 & 0 \\ 0 & p+1 & -2p & 0 \\ 0 & -2p & p+1 & 0 \\ 0 & 0 & 0 & 1-p \end{pmatrix}$$

$$\rho^\Gamma = \frac{1}{4} \begin{pmatrix} 1-p & 0 & 0 & -2p \\ 0 & p+1 & 0 & 0 \\ 0 & 0 & p+1 & 0 \\ -2p & 0 & 0 & 1-p \end{pmatrix}$$

$$\lambda_- = \frac{1}{4}(1-3p)$$



Asher Peres
1934-2005



Quantum phenomena do not occur
in a Hilbert space. They occur in a
laboratory.

— *Asher Peres* —

AZ QUOTES



Asher Peres

“Never underestimate the ingenuity of experimental physicists.”

- Asher Peres, Found. Phys. 14, 1131 (1984)

Physicist and quantum information theory pioneer, Asher Peres (1934-2005) explored connections between quantum mechanics and the theory of relativity.

#quantumchangers



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WATERLOO

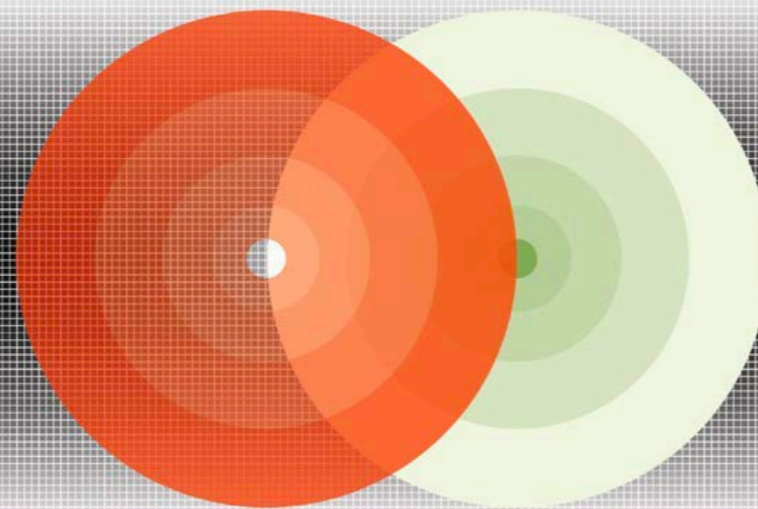
IQC

Institute for
Quantum
Computing

Quantum Theory: Concepts and Methods

by
Asher Peres


Springer Science+Business Media, LLC



Fundamental Theories of Physics

In 2×2 and 2×3 dimensions

PPT=SEP



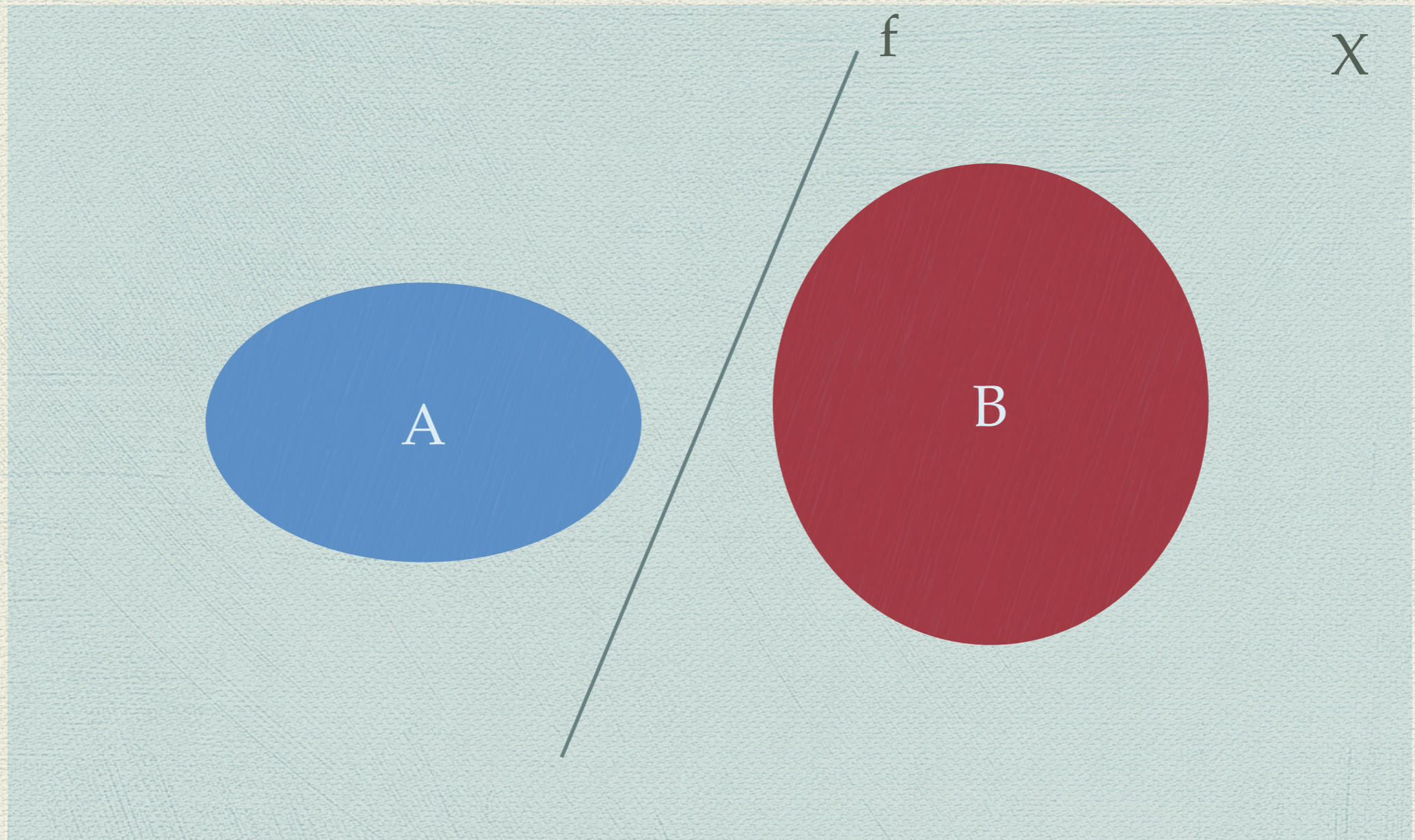
SEP=PPT

Entanglement Witnesses

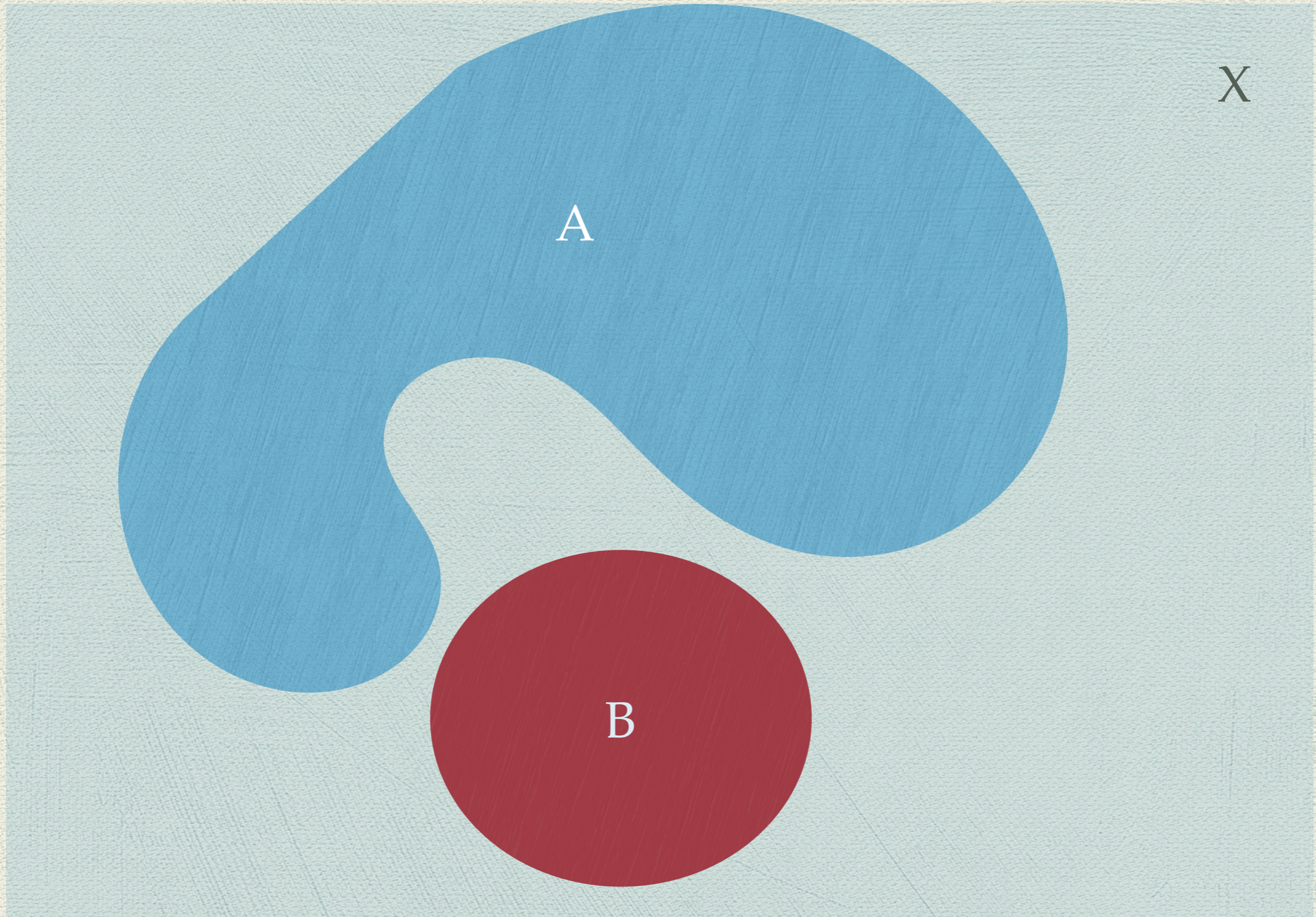
Experimental Verification of Entanglement

$Tr(W\rho)$  Statement about entanglement

Hahn-Banach Theorem

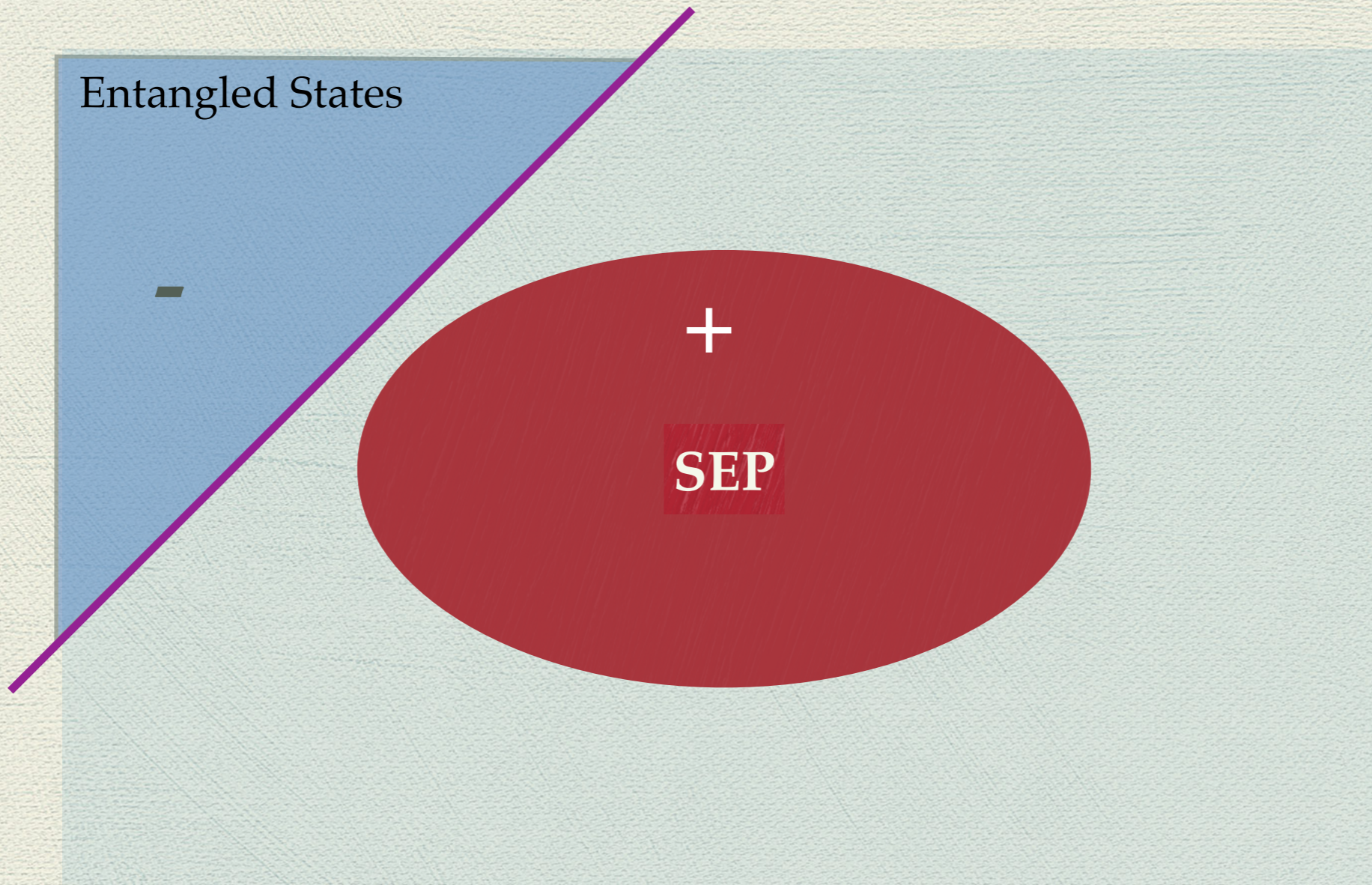


Hahn-Banach Theorem



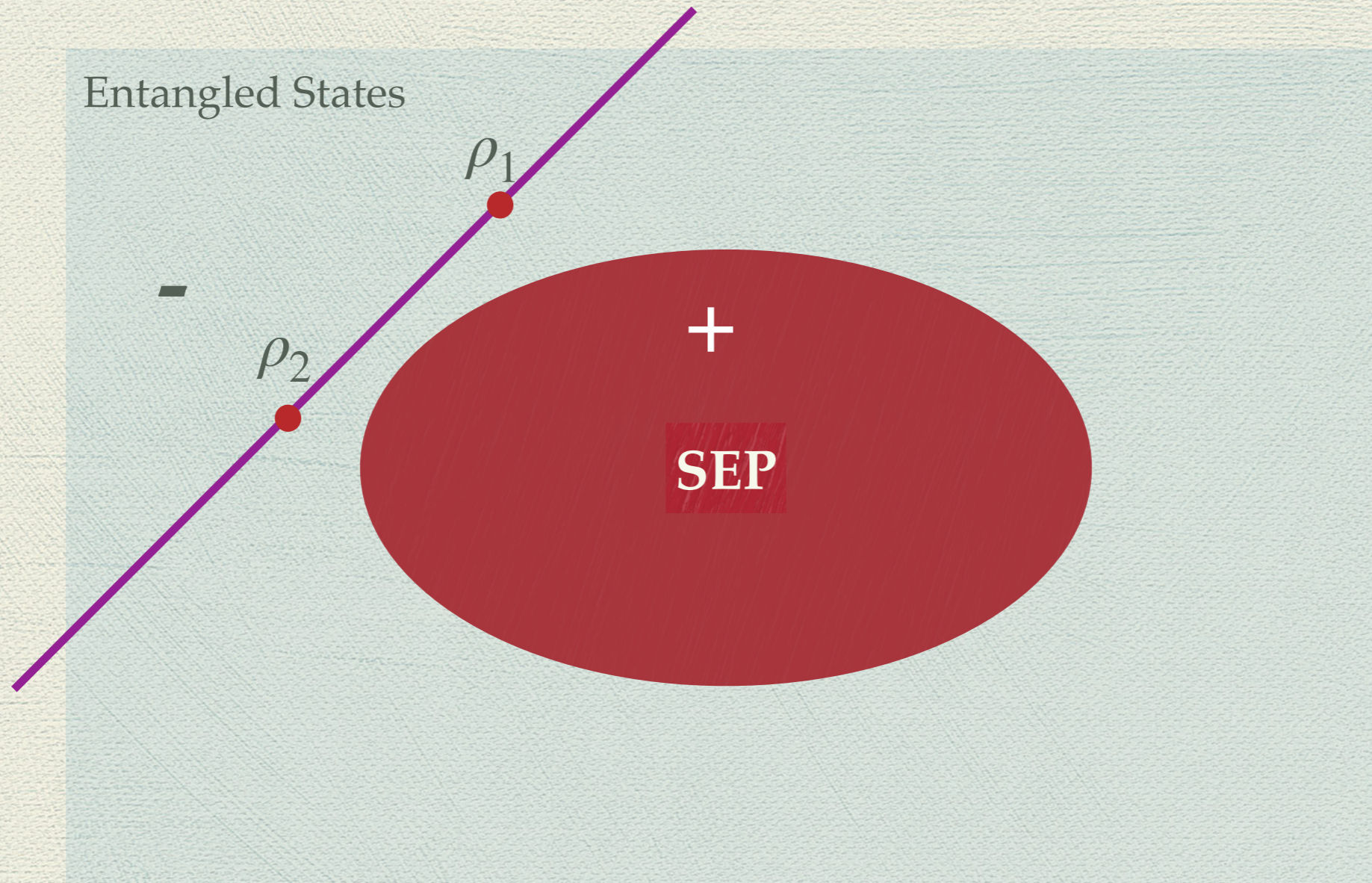
$$\text{Tr}(W\rho) = 0$$

Entangled States



$$\text{Tr}(W\rho) = 0$$

Entangled States



$$\text{Tr}(W\rho_1) = 0$$

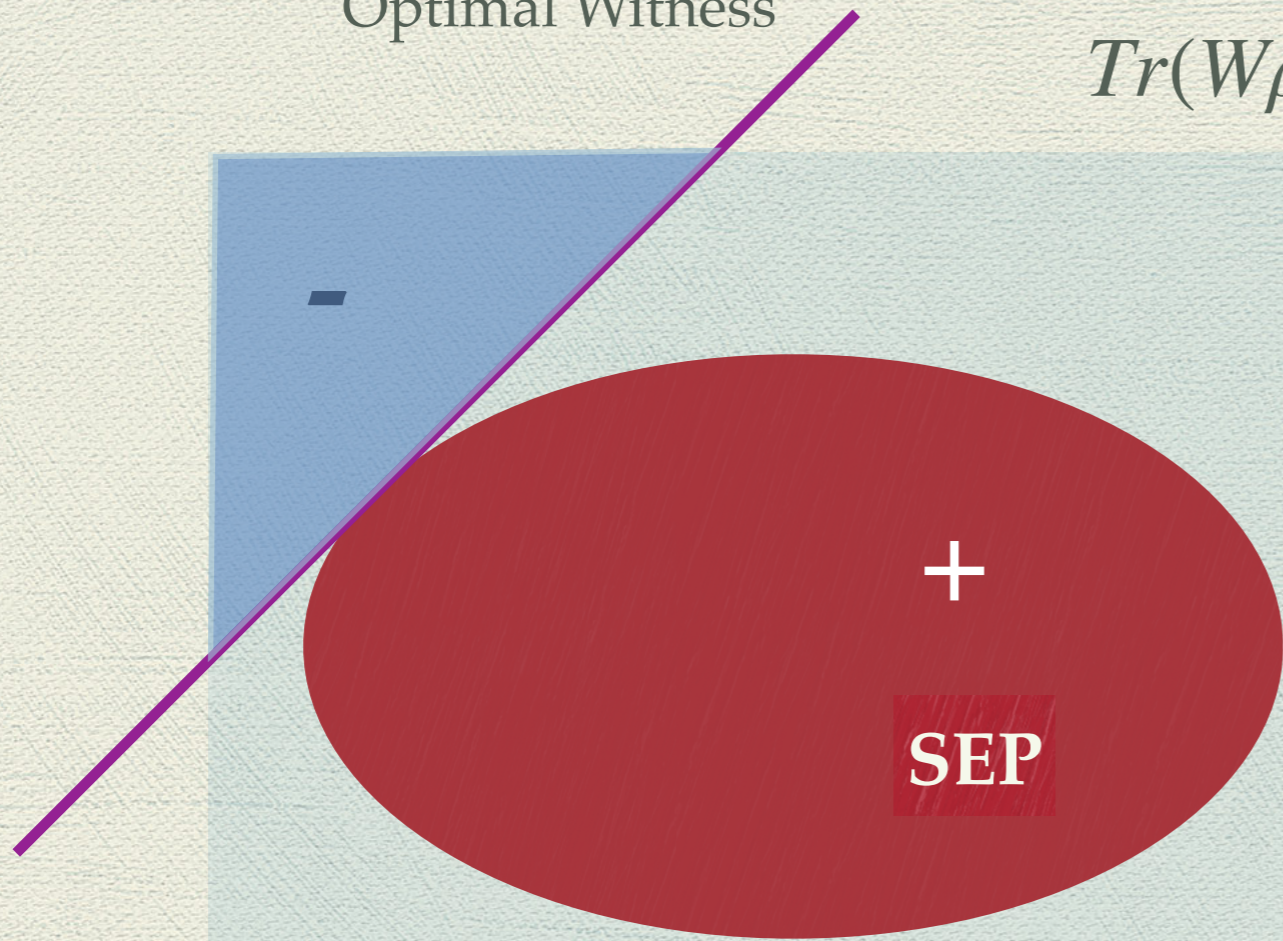
$$\text{Tr}(W\rho_2) = 0$$

$$\text{Tr}(W(\alpha\rho_1 + \beta\rho_2)) = 0$$

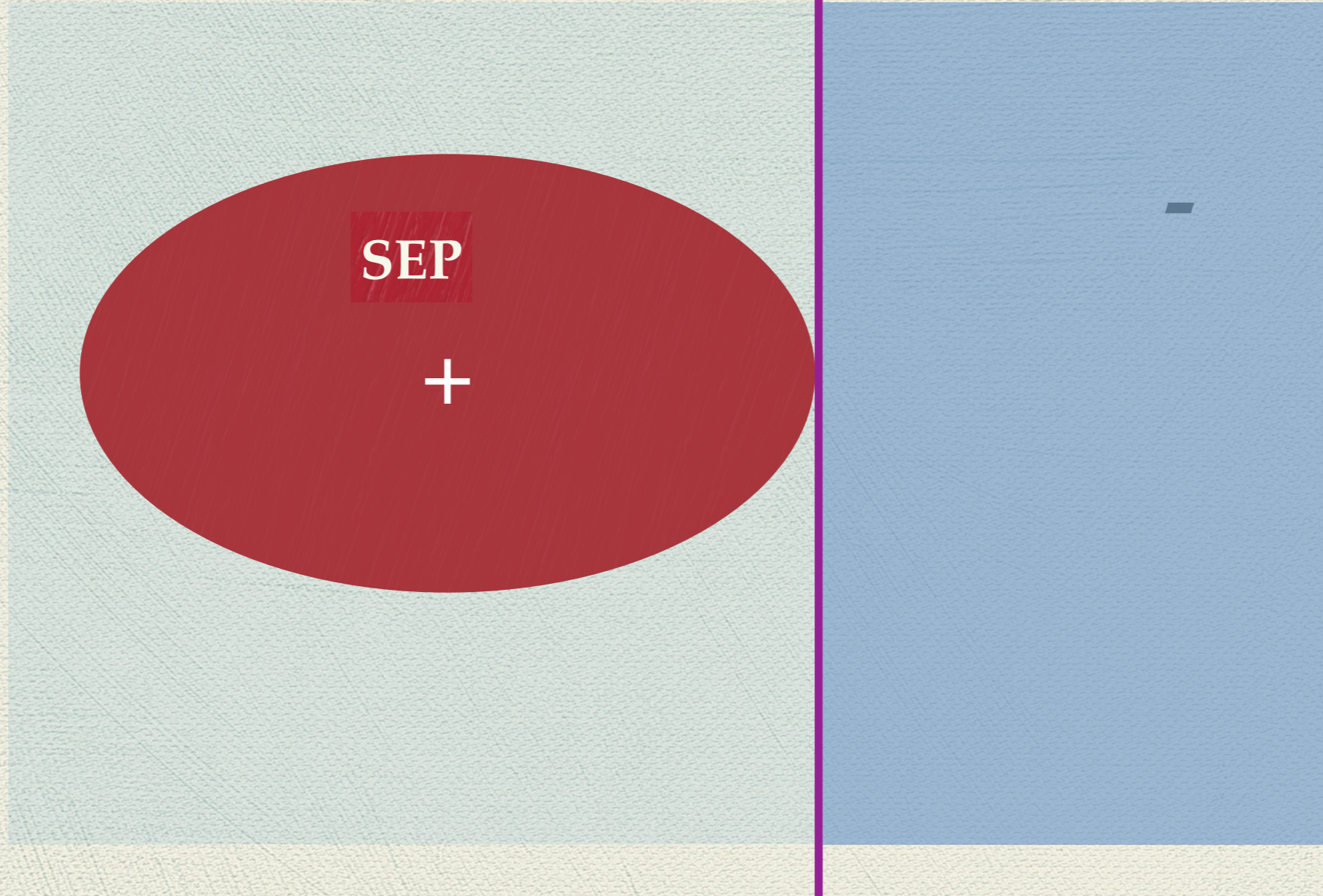
Optimal Witness

$$\text{Tr}(W\rho) = 0$$

Entangled States



$$\text{Tr}(W\rho) = 0$$



A good Witness

$$\text{Tr}(W\rho) = 0$$

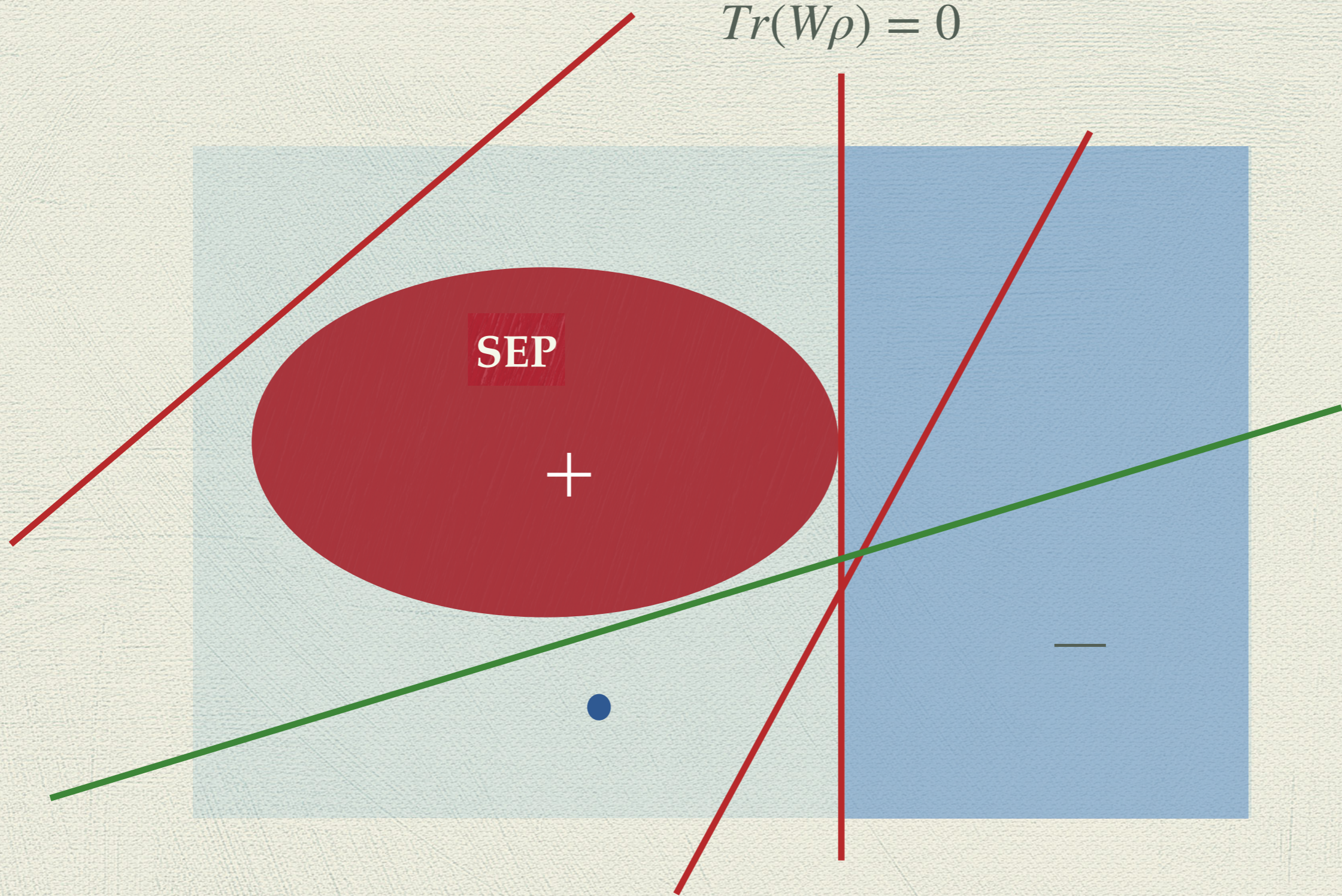
SEP

+



-

A good Witness



How to construct entanglement witnesses?

Let us learn from the PPT.

T =Transpose is a positive map

But it is not a Completely Positive Map (CPT)

$$(I \otimes T) |\phi^+\rangle\langle\phi^+| \not\geq 0$$

$$(I \otimes T) |\phi^+\rangle \langle \phi^+| \not\geq 0$$

$$W = (I \otimes T) |\phi^+\rangle \langle \phi^+|$$

$$= (I \otimes T) |i, i\rangle \langle j, j|$$

$$= |i\rangle \langle j| \otimes |j\rangle \langle i|$$

$$= |i, j\rangle \langle j, i| = P$$

$$\text{Tr}(W \rho^{sep}) > 0$$

But T is not the only positive map which is not CPT!

Let us take a map like Φ

$$W = (I \otimes \Phi) |\phi^+\rangle\langle\phi^+|$$

$$= \sum_{i,j} (I \otimes \Phi) |i,i\rangle\langle j,j|$$

$$= \sum_{i,j} |i\rangle\langle j| \otimes \Phi(|i\rangle\langle j|)$$

$$W = \sum_{i,j} |i\rangle\langle j| \otimes \Phi(|i\rangle\langle j|)$$

$$\text{Tr}[W(\rho \otimes \sigma)]$$

$$= \sum_{i,j} \text{Tr}|i\rangle\langle j|\rho \otimes \text{Tr} \Phi(|i\rangle\langle j|)\sigma$$

$$= \sum_{i,j} \rho_{j,i} \text{Tr}[\Phi(|i\rangle\langle j|)\sigma]$$

$$= \text{Tr}[\Phi(\rho^T)\sigma] \geq 0$$

Note: If A and B are positive, AB is not necessarily positive!

But

$$\text{Tr}(AB) \geq 0$$

$$\text{Tr}(AB) = \sum_{\alpha} \langle \alpha | B | \alpha \rangle \geq 0$$

To construct Witnesses,
find Positive Maps Φ which are not CTP

Then your witness will be

$$\begin{aligned} W &= (I \otimes \Phi) |\phi^+\rangle\langle\phi^+| \\ &= \sum_{i,j} |i\rangle\langle j| \otimes \Phi(|i\rangle\langle j|) \end{aligned}$$

Question: Is any witness of the form

$$W = (I \otimes \Phi) |\phi^+\rangle\langle\phi^+| ?$$

Yes:

W is a witness



$$W = (I \otimes \Phi) |\phi^+\rangle\langle\phi^+|$$

Proof:

$$W = \sum_{i,j} |i\rangle\langle j| \otimes \hat{W}_{ij}$$

$$W(|\alpha\rangle\langle\alpha| \otimes |\beta\rangle\langle\beta|) \geq 0$$

$$\alpha_j^* \alpha_i \hat{W}_{ij}(|\beta\rangle\langle\beta|) \geq 0$$

Define:

$$\Phi(|i\rangle\langle j|) := \hat{W}_{ij}$$

$$\Phi(|\alpha\rangle\langle\alpha|) := \alpha_i \alpha_j^* \hat{W}_{ij}$$

$$\langle\beta|\Phi(|\alpha\rangle\langle\alpha|)|\beta\rangle := \alpha_i \alpha_j^* \langle\beta|\hat{W}_{ij}|\beta\rangle \geq 0$$

Therefore Φ is a positive operator:

$$W = \sum_{i,j} |i\rangle\langle j| \otimes \hat{W}_{ij} = \sum_{i,j} |i\rangle\langle j| \otimes \Phi(|i\rangle\langle j|)$$

Classification of positive maps = Classification of Entanglement Witnesses

The problem:

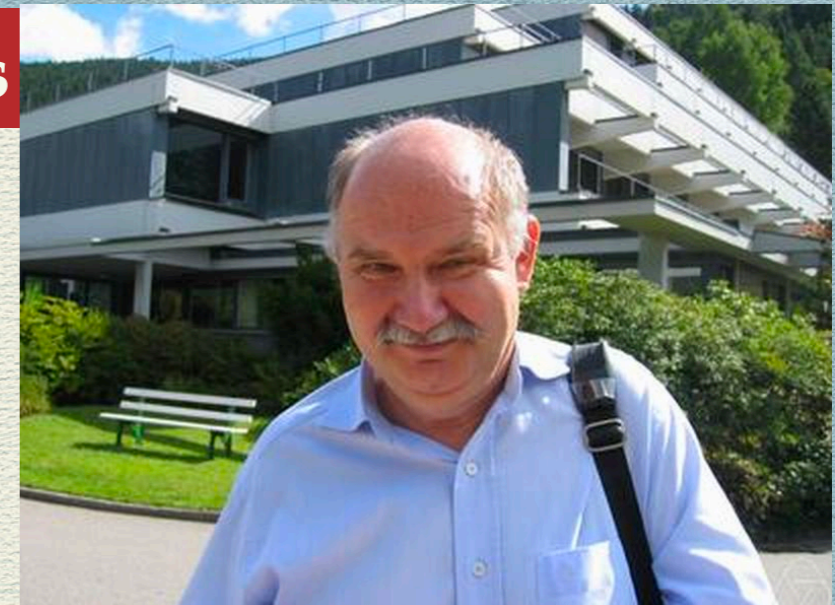
Classification of CPT's is easy.

But classification of positive maps is difficult,

A basic theorem in 2×2 and 2×3 dimensions

Only if: $\Phi : L(H_2) \longrightarrow L(H_2)$

$\Phi : L(H_2) \longrightarrow L(H_3)$



S.L. Woronowics (1976)

Every Positive maps is of the form: $\Phi = \mathcal{E}_1 + \mathcal{E}_2 \circ \mathcal{T}$

Where \mathcal{T} is the transpose map.

Decomposable maps and their corresponding witnesses.

$$\mathcal{E} = \mathcal{E}_1 + \mathcal{E}_2 \circ \mathcal{T} \quad \longrightarrow \quad W = P + (I \otimes T)Q$$

$$W = P + Q^\Gamma$$

Proof:

$$\mathcal{E} = \mathcal{E}_1 + \mathcal{E}_2 \circ \mathcal{T}$$

$$W = \left[I \otimes \mathcal{E}_1 + I \otimes \mathcal{E}_2 \circ \mathcal{T} \right] (|\phi^+\rangle\langle\phi^+|)$$

$$\left[I \otimes \mathcal{E}_1 \right] (|\phi^+\rangle\langle\phi^+|) = P$$

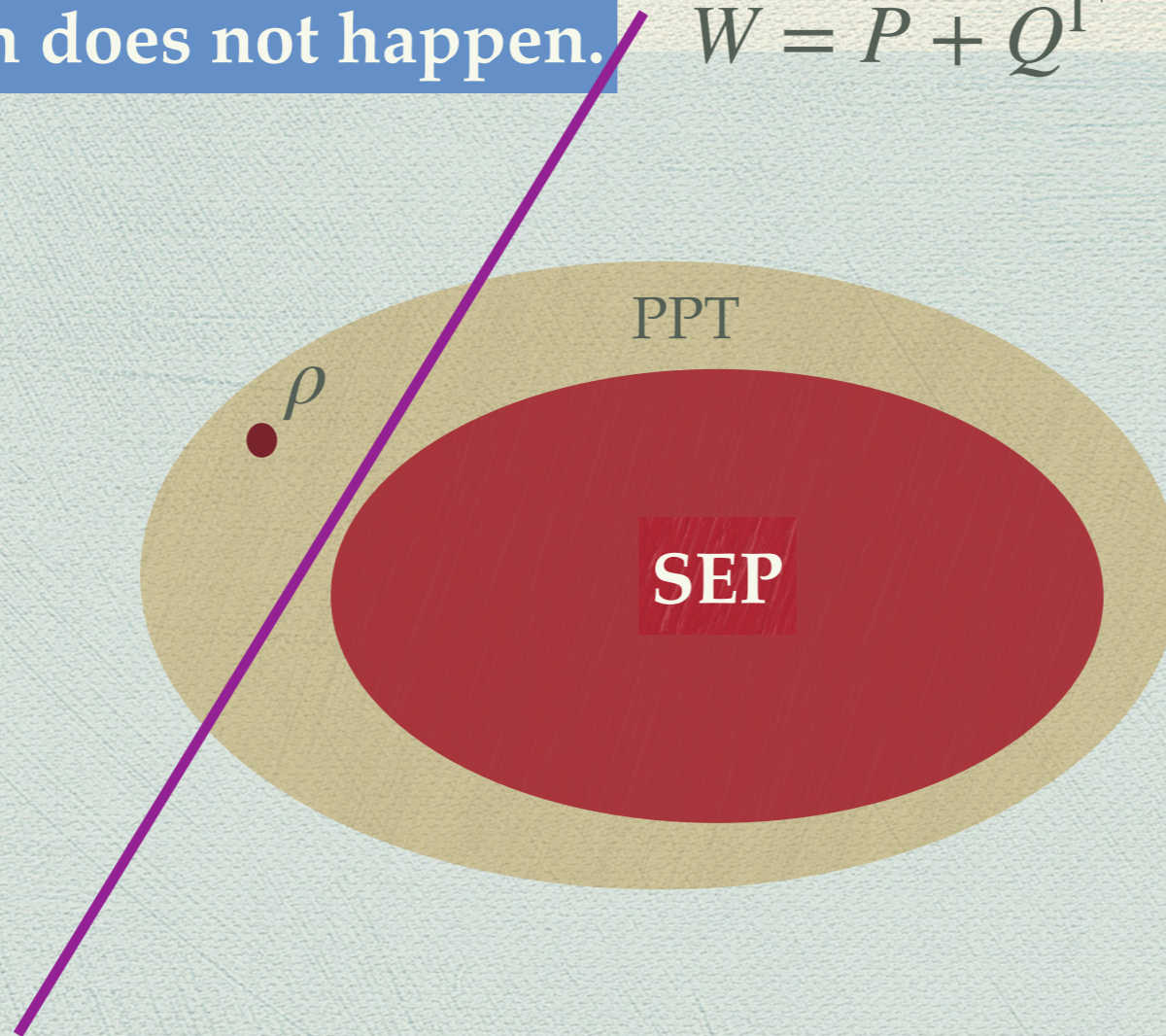
$$\left[I \otimes \mathcal{E}_2 \circ \mathcal{T} \right] (|\phi^+\rangle\langle\phi^+|) = (I \otimes T)Q$$

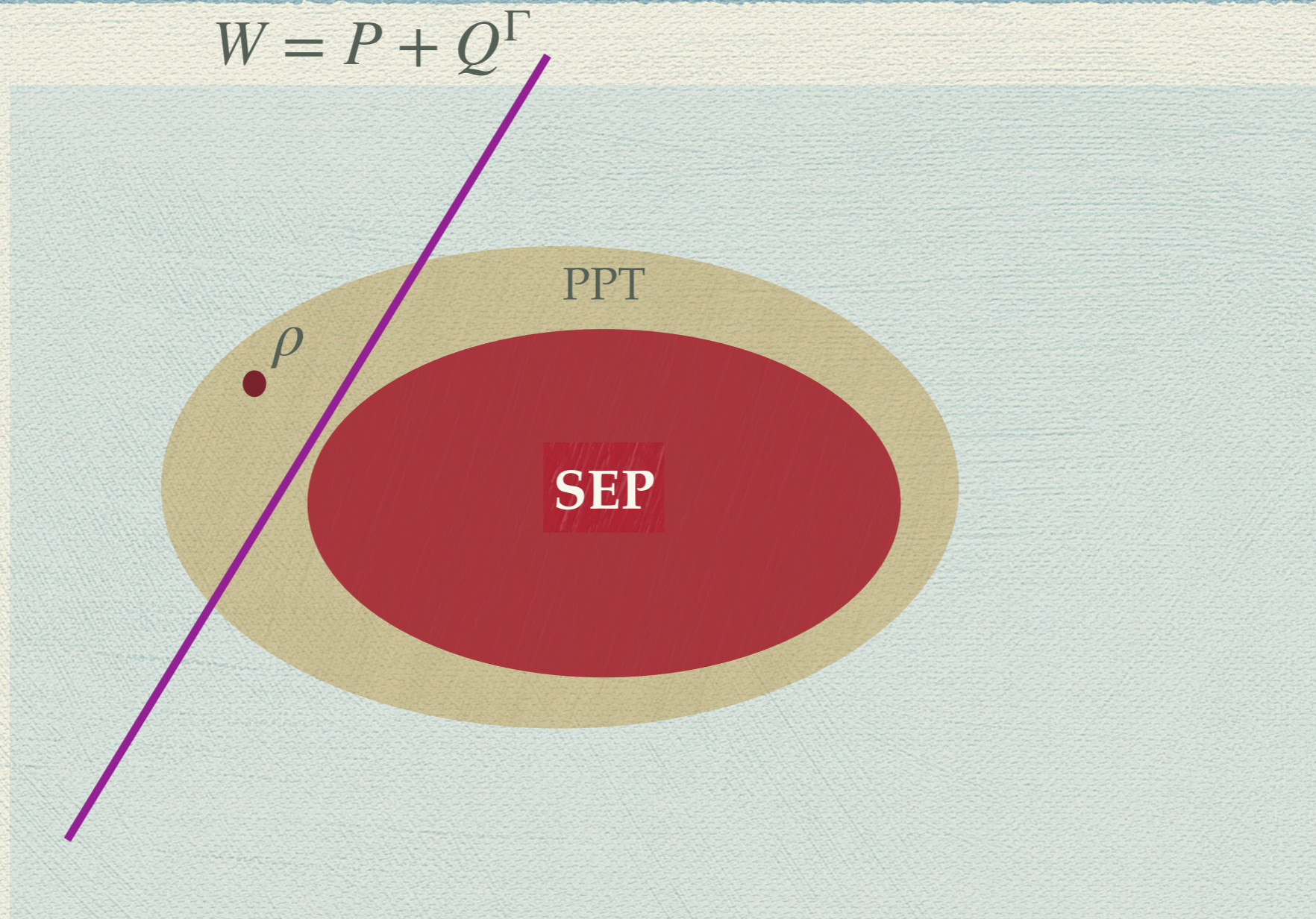
Why decomposable maps are important?

Because they detect PPT states.

This situation does not happen.

$$W = P + Q^\Gamma$$

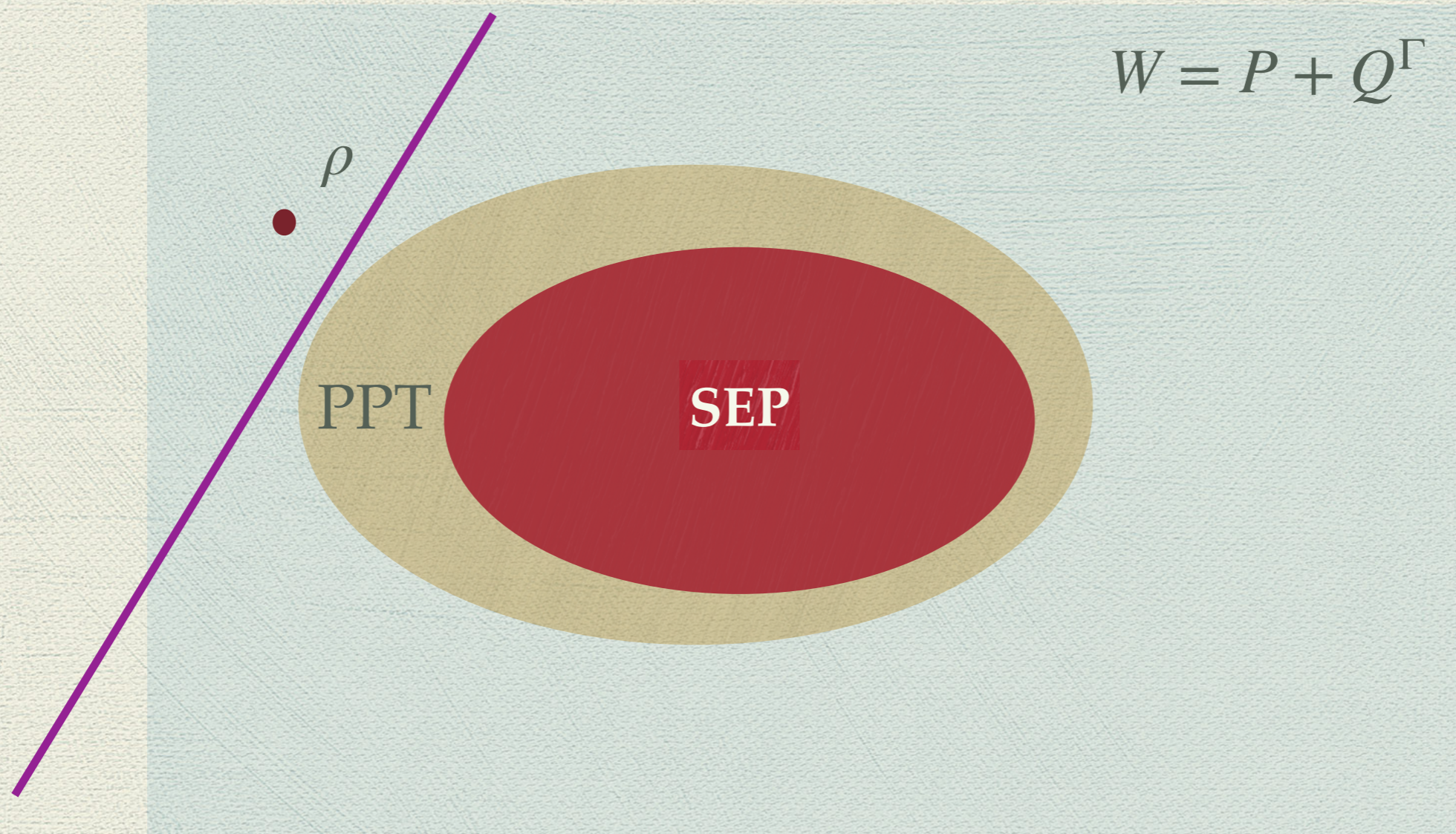




We ask if this is possible $Tr(W\rho) < 0$?

But $Tr(W\rho) = Tr(P\rho) + Tr(Q^\Gamma\rho)$ $Tr(Q^\Gamma\rho) = Tr(Q\rho^\Gamma) > 0$

This situation happens.



So if $Tr(W\rho) < 0 \longrightarrow \rho$ is not PPT!

This is the proof that :

In 2×2 and 2×3 dimensions

$PPT=SEP$



Horodecki Family

SEP=PPT

What about higher dimensions?

A journey to the 19th Century and basic mathematics.

Every Positive quadratic form is the sum of squares

$$a^2 + b^2 + 2ab = (a + b)^2$$

$$2a^2 + 5b^2 - 2ab = (a + b)^2 + (a - 2b)^2$$

$$Q_2(\mathbf{x}) = \sum_{\alpha} |Q_1^{\alpha}(\mathbf{x})|^2$$

$$Q_2(\mathbf{x}) = \sum_{\alpha} |Q_1^{\alpha}(\mathbf{x})|^2$$



$$Q_4(\mathbf{x}) = \sum_{\alpha} |Q_2^{\alpha}(\mathbf{x})|^2 \quad ?$$

$$Q_6(\mathbf{x}) = \sum_{\alpha} |Q_3^{\alpha}(\mathbf{x})|^2 \quad ?$$

David Hilbert: (1888)

Every quartic polynomial of 3 variables

$$Q_4(x, y, z) = \sum_{\alpha=1}^3 [Q_2^{(\alpha)}(x, y, z)]^2$$



David Hilbert (1862-1943)

**Hilbert (1888): For all other combinations of
($n > 3$, d) the answer is negative.**

Number of variables

Degree of polynomials

Quadratic forms and operators

$$Q(\mathbf{x}) = \sum_{i,j} \bar{x}_i Q_{ij} x_j$$

$$\hat{Q} : H_n \longrightarrow H_n$$

$$|y_i\rangle = \sum_j \hat{Q}_{ij} |x_j\rangle$$

If the quadratic form $Q(\mathbf{x})$ is positive, i.e. If $\langle \mathbf{x} | \hat{Q} | \mathbf{x} \rangle \geq 0$

We can rewrite $Q(\mathbf{x})$ in the following form:

$$Q(\mathbf{x}) = \sum_{\alpha} |(\mathbf{b}_{\alpha} \cdot \mathbf{x})|^2$$

$$\langle \mathbf{x} | \hat{Q} | \mathbf{x} \rangle = \sum_{\alpha} \langle \mathbf{x} | \mathbf{b}_{\alpha} \rangle \langle \mathbf{b}_{\alpha} | \mathbf{x} \rangle$$

$$\longrightarrow \hat{Q} = \sum_{\alpha} |\mathbf{b}_{\alpha}\rangle \langle \mathbf{b}_{\alpha}|$$

A positive linear operator corresponds to this quadratic form.

Bi-quadratic Polynomials

$$\Phi(\mathbf{x}, \mathbf{y}) = \sum_{i,j,k,l} \bar{y}_i y_j \Phi_{ij;kl} x_k \bar{x}_l$$

$$x^2 y^2 = (xy)^2$$

$$a^2 x^2 + 9b^2 y^2 + a^2 y^2 + 4b^2 x^2 - 2abxy = (ax - 3by)^2 + (ay + 2bx)^2$$



Bi-linear

Is this general?

Bi-quadratic forms and Super-operators (maps).

$$\Phi(\mathbf{x}, \mathbf{y}) = \sum_{i,j;k,l} \bar{y}_i y_j \Phi_{ij;kl} x_k \bar{x}_l$$

$$\hat{\Phi} : L(H_n) \longrightarrow L(H_n)$$

$$\hat{\Phi} : |x_k\rangle\langle x_l| \longrightarrow \sum_{i,j} \Phi_{ij;kl} |x_k\rangle\langle x_l|$$

$$\hat{\Phi} : \rho_{kl} \longrightarrow \sum_{i,j} \Phi_{ij;kl} \rho_{kl}$$

If $\Phi(\mathbf{x}, \mathbf{y})$ is a positive bi-quadratic form

→ $\hat{\Phi}$ is a positive map.

Proof:

Let ρ be a positive matrix $\rho_{kl} = \sum_{\alpha} x_k^{\alpha} \bar{x}_l^{\alpha}$

$$\begin{aligned} \langle \mathbf{y} | \hat{\Phi}(\rho) | \mathbf{y} \rangle &= \sum_{i,j} \bar{y}_i \hat{\Phi}(\rho)_{ij} y_j \\ &= \sum_{\alpha} \sum_{i,j} \bar{y}_i y_j \Phi_{ij,kl} x_k^{\alpha} \bar{x}_l^{\alpha} \end{aligned}$$

**Question: What corresponds to bi-quadratic forms
Which is a sum of squares of bi-linear forms?**

$$\Phi(\mathbf{x}, \mathbf{y}) = \sum_{\alpha} |L^{\alpha}(\mathbf{x}, \mathbf{y})|^2$$

$$L(\mathbf{x}, \mathbf{y}) = \sum_{i,l} \bar{y}_i L_{il} x_l \quad ?$$

$$\Phi(\mathbf{x}, \mathbf{y}) = \sum_{\alpha} \sum_{i,j,k,l} \bar{y}_i L_{ik}^{\alpha} x_k \quad y_j \bar{L}_{jl}^{\alpha} \bar{x}_l$$

$$\Phi_{ijkl} = \sum_{\alpha} L_{ik}^{\alpha} \bar{L}_{jl}^{\alpha}$$

$$\hat{\Phi} : \rho_{kl} \longrightarrow \sum_{i,j} \Phi_{ij;kl} \rho_{kl}$$

$$\hat{\Phi} : \rho_{kl} \longrightarrow \sum_{i,j} \sum_{\alpha} L_{ik}^{\alpha} \bar{L}_{jl}^{\alpha} \rho_{kl}$$

$$\hat{\Phi} : \rho \longrightarrow \sum_{\alpha} L^{\alpha} \rho L^{\alpha \dagger}$$

Therefore if every positive bi-quadratic form
is the sum of squares of bi-linear forms

Then every positive map is a CPT.

**Question: What corresponds to bi-quadratic forms
Which is a sum of squares of bi-linear forms?**

Choi (1975):

The answer is NO.

Positive Semidefinite Biquadratic Forms

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Berkeley, California 94720*

Submitted by Chandler Davis



Counter-Example:

$$\begin{aligned} F = & x_1^2 y_1^2 + x_2^2 y_2^2 + x_3^2 y_3^2 \\ & - 2(x_1 x_2 y_1 y_2 + x_2 x_3 y_2 y_3 + x_3 x_1 y_3 y_1) \\ & + 2(x_1^2 y_1^2 + x_2^2 y_2^2 + x_3^2 y_3^2) \end{aligned}$$



So every positive map Φ is not a CPT \mathcal{E} .

What about

$$\Phi = \mathcal{E}_1 + \mathcal{E}_2 \circ \mathcal{T}$$

Again the answer is NO!

$$\text{If } \hat{\Phi} = \mathcal{E}_1 + \mathcal{E}_2 \circ \mathcal{T}$$

$$\hat{\Phi}(\rho) = \sum_{\alpha} F_{\alpha} \rho F_{\alpha}^{\dagger} + \sum_{\beta} G_{\beta} \rho^T G_{\beta}^{\dagger}$$

$$\Phi_{ij,kl} \rho_{kl} = \sum_{\alpha} (F_{\alpha})_{ik} \rho_{kl} (F_{\alpha}^{\dagger})_{lj} + \sum_{\beta} (G_{\beta})_{ik} \rho_{lk} (G_{\beta}^{\dagger})_{lj}$$

$$\Phi_{ij,kl} = \sum_{\alpha} (F_{\alpha})_{ik} (F_{\alpha}^{\dagger})_{lj} + \sum_{\beta} (G_{\beta})_{il} (G_{\beta}^{\dagger})_{kj}$$

$$\Phi_{ij,kl} = \sum_{\alpha} (F_{\alpha})_{ik} (F_{\alpha}^{\dagger})_{lj} + \sum_{\beta} (G_{\beta})_{il} (G_{\beta}^{\dagger})_{kj}$$

$$\Phi_{ij,kl} = \sum_{\alpha} (F_{\alpha})_{ik} \overline{(F_{\alpha})_{jl}} + \sum_{\beta} (G_{\beta})_{il} \overline{(G_{\beta})_{jk}}$$

$$\begin{aligned} \sum_{ijkl} y_i \bar{y}_j \Phi_{ij,kl} x_k \bar{x}_l &= \sum_{\alpha} y_i \bar{y}_j (F_{\alpha})_{ik} \overline{(F_{\alpha})_{jl}} x_k \bar{x}_l \\ &+ \sum_{\beta} y_i \bar{y}_j (G_{\beta})_{il} \overline{(G_{\beta})_{jk}} x_k \bar{x}_l \\ &= \sum_{\alpha} |(\mathbf{y}^T F^{\alpha} \mathbf{x})|^2 \\ &+ \sum_{\beta} |(\mathbf{y}^T G^{\beta} \bar{\mathbf{x}})|^2 \end{aligned}$$

Therefore if

$$\hat{\Phi} = \mathcal{E}_1 + \mathcal{E}_2 \circ \mathcal{T}$$

We conclude that

$$\Phi(\mathbf{x}, \mathbf{y}) = \sum_{\alpha} |(\mathbf{y}^T F^{\alpha} \mathbf{x})|^2 + \sum_{\beta} |(\mathbf{y}^T G^{\beta} \bar{\mathbf{x}})|^2$$

which we know is not always true.

Summary:

For qubits-qubit and qubit-qutrit, any Witness is of the form

$$W = P + Q^\Gamma$$

In higher dimensions, there are other forms of witnesses.

End of Part I